

QDs-Cavity Approach to Controlled Quantum Teleportation of GHZ-Like State

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Abstract

The experimental scheme for controlled quantum teleportation of tripartite GHZ-like state is presented. With the entanglement generating through the interaction between the quantum dots in microcavities and a single photon, the controlled teleportation can be realized by virtue of Faraday rotation, single photon detection and electron spin orientation measurement. The success probability of the scheme can reach 1 if the cavities are switchable to choice the appropriate Faraday rotation angle. The scheme can be easily generalized to the teleportation of multipartite GHZ-like state.

Keywords: Controlled quantum teleportation; QDs-cavity system; Faraday rotation; Multipartite GHZ-Like state

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1. Introduction

Quantum information physics has developed into an extensive field involving the subjects of physics, informatics, computer, engineering science, chemistry and mathematics *etc.* and the corresponding research work of the past several years has greatly clarified both the theoretical potential and the experimental challenges of quantum information physics.

To produce working laboratory devices that perform this profoundly new form of information processing is one of main subjects of quantum information physics. By now, many experimental realizations have been performed with optical systems [1-2], NMR techniques [3], super-conducting device physics [4], electron physics [5], quantum dot [6-10], ion-trap [11] and cavity-QED technique [12] and so on, but today's implementation of quantum computer still meets serious challenges, such as quantum decoherence and quantum measurement, therefore many efforts have been made to design academic protocols based on different working laboratory devices.

Since the seminal work of Loss and Divincenzo [7] and papers of Obermayer, Teich and Mahler [13], there has been increasing interest in quantum information processing (QIP) with quantum dots (QDs), which are propitious for integration and miniaturization. Subsequently, several theoretical proposals [14-16] for quantum teleportation (QT) of electron spin in a QD have been proposed due to the relative robustness of the electron spin against decoherence. Furthermore, with QDs placed in a microcavity and illuminated by laser beams, and with interaction between QDs mediated by a single-mode cavity field, the scheme of Imamoglu et al. [9] realizes controlled interactions between two distant QDs to a certain extent. We call such composite system as QDs-cavity system. It is worth noting that a crucial condition for a realization of QIP in such models is a strong coupling of a single QD to a single mode of microcavity (or nanocavity) of a high quality factor (high-Q). Quite recently, such random couplings [19-20] or even deterministic couplings [21] have been observed experimentally.

Considering the security of quantum communication and the expansibility of quantum network, we focus on the controlled quantum teleportation (CQT) of tripartite GHZ-like state. CQT scheme was firstly proposed in 1998 [22]. The feature of this scheme is that the teleportation between the sender and the receiver is under the control of the third party (the controller). In this article, the teleportation scheme is based on QDs coupled through a microcavity mode and on the controlling nanostructure system composed of microcavity with QDs inside. The entanglement in microcavity is generated by the interaction between the electron spin and photon. Using single photon measurement [17, 23-24], Faraday rotation [17, 23] and electron spin orientation measurement, the CQT can be implemented perfectly and we find that the success probability of the scheme can reach 1 if the cavities are switchable to choose the appropriate Faraday rotation angle.

The article is organized as follows: we first briefly review the generation of the entangled state by Faraday rotation in QDs-cavity system, then propose the controlled teleportation scheme of tripartite GHZ-like state,

which is under the control of a three-quantum-dot system.

2. Generation of the entangled state by Faraday rotation in QD-cavity system

When a photon interacts with an electron spin of QD in microcavity, the photon will excite two virtual processes in the microcavity and Faraday rotation will happen in these two virtual processes [17].

The virtual processes in Figure 1 can be described as follow: For excess electron's spin- \downarrow , light of $\sigma_{(z)}^-$ polarization excites an electron and a heavy hole (upper right), and light of $\sigma_{(z)}^+$ polarization excites an electron and a light hole (lower left). The final occupations of electron states in the dot are the same as the initial occupations, but a phase shift $e^{iS_0^{hh}}$ of $\sigma_{(z)}^-$ or $e^{iS_0^{lh}}$ of $\sigma_{(z)}^+$ will appear. Here $S_0^{hh} = \Omega_{hh} T$ and $S_0^{lh} = \Omega_{lh} T$ are for heavy and light hole respectively, $\Omega_{hh} = V_{hh}^2 / \hbar^2 \omega_d = 3V_{hh}^2 / \hbar^2 \omega_d$ and $\Omega_{lh} = V_{lh}^2 / \hbar^2 \omega_d$ are the rates of Faraday rotation for

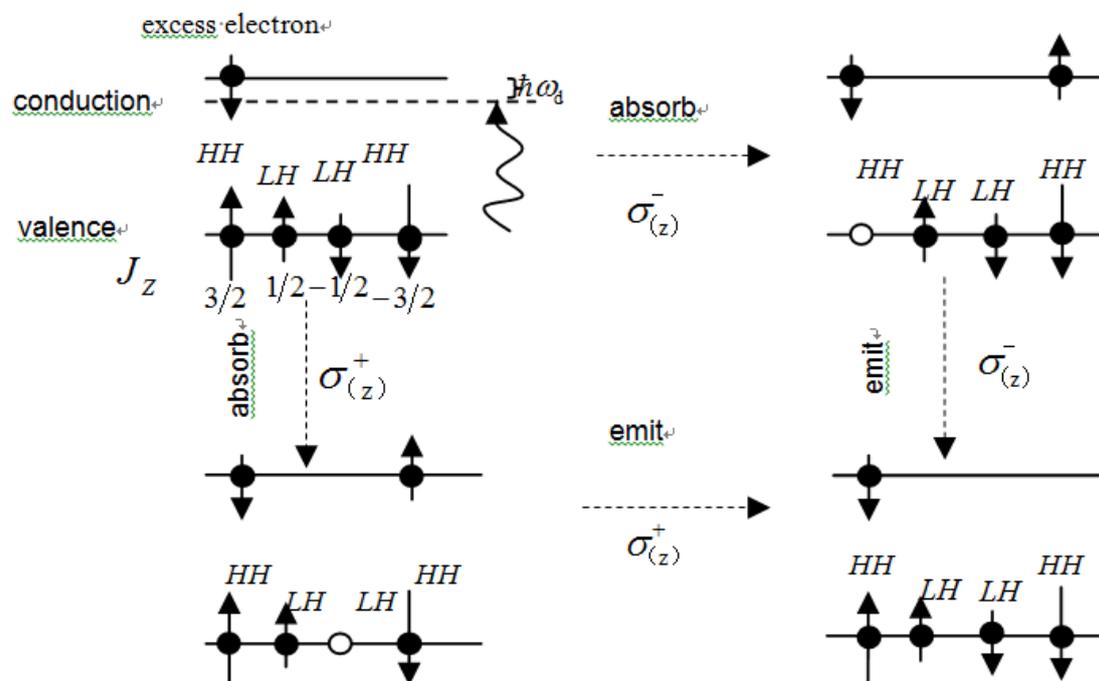


Figure 1. The virtual process will happen after the spin entangles with the photon in microcavity. For spin- \downarrow , light of $\sigma_{(z)}^-$ polarization excites an electron and a heavy hole (upper right), and light of $\sigma_{(z)}^+$ polarization excites an electron and a light hole (lower left). J_z denotes angular momentum quantum number.

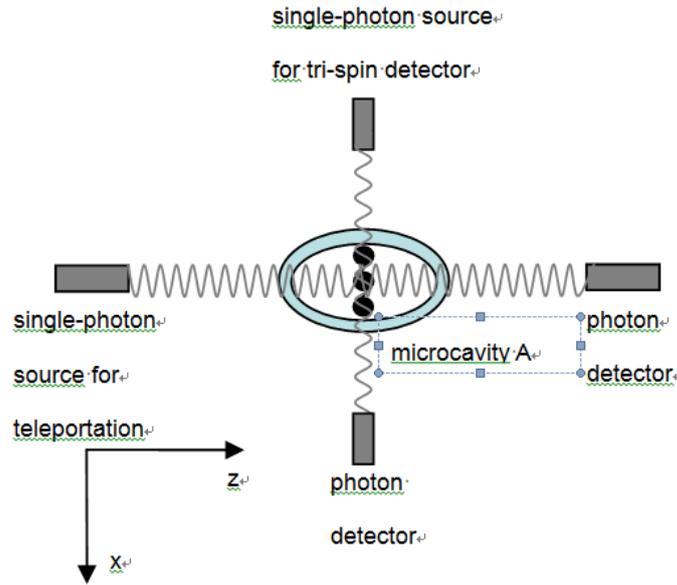


Figure 2. The generation of entanglement between electron spins and photon by Faraday rotation in QD-cavity system. Spin detection is performed by photon propagating along x.

these virtual transitions, $\hbar\omega_d$ is the detuning of the non-resonant photon, V_{hh} and V_{lh} are the transition matrix elements for exciting an electron and a heavy hole or an electron and a light hole in the Faraday rotation process and $V_{hh}, V_{lh} \ll \hbar\omega_d$, T is the residence time of photon in microcavity. Therefore, under the basis $\{|\downarrow\rangle|\sigma_{(z)}^-\rangle, |\downarrow\rangle|\sigma_{(z)}^+\rangle, |\uparrow\rangle|\sigma_{(z)}^-\rangle, |\uparrow\rangle|\sigma_{(z)}^+\rangle\}$, the full Hamiltonian of these processes in the rotation frame can be written as

$$H = \begin{pmatrix} E_c - \hbar\omega_d & V_{hh} & 0 & 0 \\ V_{lh} & E_c & 0 & 0 \\ 0 & 0 & E_c - \hbar\omega_d & V_{hh} \\ 0 & 0 & V_{lh} & E_c \end{pmatrix}, \quad (1)$$

where E_c is the energy of the excess electron.

Similarly, if excess electron's spin is \uparrow , the corresponding Hamiltonian of these processes under the basis $\{|\uparrow\rangle|\sigma_{(z)}^+\rangle, |\uparrow\rangle|\sigma_{(z)}^-\rangle, |\downarrow\rangle|\sigma_{(z)}^+\rangle, |\downarrow\rangle|\sigma_{(z)}^-\rangle\}$ can be written as

$$H = \begin{pmatrix} E_c - \hbar\omega_d & V_{hh} & 0 & 0 \\ V_{hh} & E_c & 0 & 0 \\ 0 & 0 & E_c - \hbar\omega_d & V_{lh} \\ 0 & 0 & V_{lh} & E_c \end{pmatrix}. \quad (2)$$

Since the dimension of QD is on the order of nanometer, if we place three QDs in a microcavity (Figure 2), it is possible that the total size of QD array is on the scale of nanometer and the wavelength of photon of several hundreds of nanometers is larger than the dimension of three-QD system. Therefore, when the photon propagates into the microcavity, the single photon can interact with the cluster of three QDs simultaneous-

ly [23]. Meanwhile, the interaction produces a conditional single photon Faraday rotation [17, 23]. Now, we describe the Faraday rotation angle quantitatively and the generation of entanglement between electron spins and photon in QD-cavity system in detail. Suppose the original state of electron spins in QD system A is prepared as $|\phi_A\rangle = |\uparrow\uparrow\uparrow\rangle_A$ with quantization along z-direction and the photon propagates along the opposite of z-direction with linear polarization initially along the x-direction, i.e. the spin-photon wave-function before interaction can be written as

$$\begin{aligned} |\psi_{Ap}(0)\rangle &= |\leftrightarrow\rangle_p |\phi_A\rangle = (|\sigma_{(z)}^+\rangle_p + |\sigma_{(z)}^-\rangle_p) |\phi_A\rangle / \sqrt{2} \\ &= |\psi_{Ap}^+(0)\rangle + |\psi_{Ap}^-(0)\rangle. \end{aligned} \quad (3)$$

In order to distinguish clearly each step of the teleportation, we introduce the times 0 and T . The evolution state of the above system is in the Hilbert space with basis

$$\begin{aligned} &\{|\uparrow\uparrow\uparrow\rangle|\sigma_{(z)}^+\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_1\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_2\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_3\rangle; \\ &|\uparrow\uparrow\uparrow\rangle|\sigma_{(z)}^-\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_1\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_2\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_3\rangle\}, \end{aligned} \quad (4)$$

here, $(\downarrow, hh)_i$ and $(\downarrow, lh)_i$ express the electron-hole pair state excited inside of QD i . The Hilbert space above is the direct sum of two subspaces: $\{|\uparrow\uparrow\uparrow\rangle|\sigma_{(z)}^+\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_1\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_2\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_3\rangle\}$ and $\{|\uparrow\uparrow\uparrow\rangle|\sigma_{(z)}^-\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_1\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_2\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_3\rangle\}$. Either subspace under the evolution is closed.

Considering the characteristic of orthogonality and linearity, we analyze two subspaces separately. In subspace $\{|\uparrow\uparrow\uparrow\rangle|\sigma_{(z)}^+\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_1\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_2\rangle, |\uparrow\uparrow\uparrow(\downarrow, hh)_3\rangle\}$, the corresponding Hamiltonian is

$$H_+ = \begin{pmatrix} E_c - \hbar\omega_d & V_{hh} & V_{hh} & V_{hh} \\ V_{hh} & E_c & 0 & 0 \\ V_{hh} & 0 & E_c & 0 \\ V_{hh} & 0 & 0 & E_c \end{pmatrix} \quad (5)$$

After the interaction between the initially unentangled photon with the spins in the three-QD system, the evolution state of the time T becomes

$$H_- = \begin{pmatrix} E_c - \hbar\omega_d & V_{lh} & V_{lh} & V_{lh} \\ V_{lh} & E_c & 0 & 0 \\ V_{lh} & 0 & E_c & 0 \\ V_{lh} & 0 & 0 & E_c \end{pmatrix} \quad (7)$$

and the evolution state of the time T is

$$\begin{aligned} |\psi_{Ap}^-(T)\rangle &= e^{-\frac{i}{\hbar}HT} |\psi_{Ap}^-(0)\rangle = e^{-\frac{i}{\hbar}HT} |\psi_{Ap}^-(0)\rangle \\ &= \frac{1}{\sqrt{2}} e^{-\frac{i}{\hbar}(E_c - \hbar\omega_d)T} e^{\frac{i3V_{lh}^2}{\hbar\hbar\omega_d}T} |\uparrow\uparrow\uparrow\rangle_e |\sigma_{(z)}^-\rangle_p. \end{aligned} \quad (8)$$

So the evolution state of the time T in the full space can be written as

$$\begin{aligned} |\psi_{Ap}^-(T)\rangle &= |\psi_{Ap}^+(T)\rangle + |\psi_{Ap}^-(T)\rangle \\ &= e^{-\frac{i}{\hbar}(E_c - \hbar\omega_d)T} e^{\frac{i3S_0^{hh} + 3S_0^{lh}}{2}T} \{ |-3S_0/2\rangle |\uparrow\uparrow\uparrow\rangle_e \}. \end{aligned} \quad (9)$$

where $S_0 = S_0^{hh} - S_0^{lh}$, $|\varphi\rangle = (e^{-i\varphi} |\sigma_{(z)}^+\rangle + e^{i\varphi} |\sigma_{(z)}^-\rangle) / \sqrt{2}$ and S_0^{hh} , S_0^{lh} have been defined previously.

Similarly, if the photon which is initially linearly polarized in the x-direction interacts with three QDs, which are prepared in $|\varphi_A^-\rangle = |\downarrow\downarrow\downarrow\rangle_A$, we can get the evolution state of the time T as

$$\begin{aligned} |\psi_{Ap}^-(T)\rangle &= \frac{1}{\sqrt{2}} e^{-\frac{i}{\hbar}(E_c - \hbar\omega_d)T} \left\{ e^{\frac{i3V_{lh}^2}{\hbar\hbar\omega_d}T} |\downarrow\downarrow\downarrow\rangle_e |\sigma_{(z)}^-\rangle_p + e^{\frac{i3V_{lh}^2}{\hbar\hbar\omega_d}T} |\downarrow\downarrow\downarrow\rangle_e |\sigma_{(z)}^-\rangle_p \right\} \\ &= e^{-\frac{i}{\hbar}(E_c - \hbar\omega_d)T} e^{\frac{i3S_0^{hh} + 3S_0^{lh}}{2}T} (|+3S_0/2\rangle_p |\downarrow\downarrow\downarrow\rangle_A). \end{aligned} \quad (10)$$

When the state of original spins in QDs of system A is prepared in the state $|\varphi_A^-\rangle = a|\uparrow\uparrow\uparrow\rangle_A + b|\downarrow\downarrow\downarrow\rangle_A$, we can easily get the the evolution state of the time T ,

$$\begin{aligned} |\psi_{Ap}^-(T)\rangle &= e^{-\frac{i}{\hbar}HT} |\psi_{Ap}^-(0)\rangle = e^{-\frac{i}{\hbar}HT} (a|\uparrow\uparrow\uparrow\rangle_A + b|\downarrow\downarrow\downarrow\rangle_A) |\leftrightarrow\rangle_p \\ &= e^{-\frac{i}{\hbar}(E_c - \hbar\omega_d)T} e^{\frac{i3S_0^{hh} + 3S_0^{lh}}{2}T} a|\uparrow\uparrow\uparrow\rangle_A |-3S_0/2\rangle_p + b|\downarrow\downarrow\downarrow\rangle_A |+3S_0/2\rangle_p. \end{aligned} \quad (11)$$

This means that spin-photon interaction in the microcavity A produces a conditional single-photon Faraday rotation around the z-axis by the angle $\pm \frac{3S_0}{2}$ and the angle can be changed by using a switchable cavity [12]. If the state of electron spin is $|\uparrow\uparrow\uparrow\rangle_A$ ($|\downarrow\downarrow\downarrow\rangle_A$), the linear polarization of the incoming photon is rotated the angle $-\frac{3S_0}{2}$ ($+\frac{3S_0}{2}$). Therefore, from formula (11), we find that Faraday rotation and entanglement between the elec-

$$\begin{aligned} |\psi_{Ap}^+(T)\rangle &= e^{-\frac{i}{\hbar}HT} |\psi_{Ap}^+(0)\rangle = e^{-\frac{i}{\hbar}HT} |\psi_{Ap}^+(0)\rangle \\ &= \frac{1}{\sqrt{2}} e^{-\frac{i}{\hbar}(E_c - \hbar\omega_d)T} e^{\frac{i3V_{lh}^2}{\hbar\hbar\omega_d}T} |\uparrow\uparrow\uparrow\rangle_e |\sigma_{(z)}^+\rangle_p. \end{aligned} \quad (6)$$

The Hamiltonian in subspace $\{ |\uparrow\uparrow\uparrow\rangle |\sigma_{(z)}^-\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_1\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_2\rangle, |\uparrow\uparrow\uparrow(\downarrow, lh)_3\rangle \}$ is

tron spins and the photon are generated in QDs-cavity system, it can be used in the CQT in section III.

3. The controlled quantum teleportation scheme of tripartite GHZ-Like state in QDs-cavity system

In this section, we describe in detail the CQT procedure of tripartite GHZ-like state encoded in the electron spins in the QD system. In order to realize the teleportation, the photon entangled with the three QDs in microcavity A propagates to microcavity C and then interacts with the QD in microcavity C. Subsequently, the photon which carries quantum information of QDs in microcavity A and of controlling system C travels from C to B and then interacts with the three QDs in microcavity B. The corresponding teleportation diagram is shown in Figure 3. In microcavity A, we consider $|\varphi_A\rangle = a|\uparrow\uparrow\uparrow\rangle_A + b|\downarrow\downarrow\downarrow\rangle_A$, ($|a|^2 + |b|^2 = 1$), a tripartite GHZ-like state [25-27], encoded in the electron spins which correspond to the three-QD system, while the controlling system C is in $|\varphi_C\rangle = (|\uparrow\rangle_C + |\downarrow\rangle_C) / \sqrt{2} = |\leftrightarrow\rangle_C$, the QDs in destination system B is prepared in the other GHZ state [28-31] $|\varphi_B\rangle = (|\uparrow\uparrow\uparrow\rangle_B + |\downarrow\downarrow\downarrow\rangle_B) / \sqrt{2}$.

In order to distinguish clearly each step of the teleportation, we introduce the times $t_{A_1} < t_{A_2} < t_{C_1} < t_{C_2} < t_{B_1} < t_{B_2}$. Suppose the photon propagating in the z direction is initially linearly polarized in the x direction, so the original spin-photon wavefunction before interaction between the initial unentangled photon and the three entangled QDs in microcavity A is

$$|\psi_{pA}(t_{A_1})\rangle = |\varphi\rangle_A |\leftrightarrow\rangle_p. \quad (12)$$

After the interaction, the generalization of genuine four-particle entanglement between the single photon and the three entangled QDs encoded in the electron spin states is realized by the conditional Faraday rotation [17, 23] which depends on the photon polarization and spin orientation in QDs-cavity system A. According to section II, after neglecting overall phase factor, the evolution state of the composite system, composed of a single photon and three QDs in microcavity A, is $|\psi_{pA}(t_{A_2})\rangle = a|\uparrow\uparrow\uparrow\rangle_A |-3S_0/2\rangle_p + b|\downarrow\downarrow\downarrow\rangle_A |+3S_0/2\rangle_p$. (13)

Subsequently, the photon with quantum information of three QDs in microcavity A travels to the controlling system C and then interacts with three QDs in system C, it will generate the entanglement between the above entangled four-particle state and the three entangled QDs in microcavity C. Before interaction, the composite system, composed of the single photon, the three

entangled QDs in microcavity A and the three entangled QDs in the controlling system C, is in a state

$$|\psi_{pAC}(t_{C_1})\rangle = |\psi_{pA}(t_{A_2})\rangle |\varphi_C\rangle. \quad (14)$$

After interaction, the state of the composite system becomes

$$\begin{aligned} |\psi_{pAC}(t_{C_2})\rangle &= (a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |-2S_0\rangle_p + a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |-S_0\rangle_p \\ &\quad + b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |+S_0\rangle_p + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |+2S_0\rangle_p) / \sqrt{2} \\ &= [|+S_0\rangle_p (-a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C + b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C) \\ &\quad + |+2S_0\rangle_p (-a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C)] / \sqrt{2}. \end{aligned} \quad (15)$$

Finally, the photon with quantum information of four QDs in microcavity A and controlling system C propagates to the destination system B. Before interaction, the corresponding state of the composite system of photon, A, B and C is

$$|\psi_{pACB}(t_{B_1})\rangle = |\psi_{pAC}(t_{C_2})\rangle |\varphi_B\rangle. \quad (16)$$

After interaction, the state of the composite system becomes

$$\begin{aligned} |\psi_{pACB}(t_{B_2})\rangle &= (a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |-7S_0/2\rangle_p + a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |-S_0/2\rangle_p \\ &\quad + a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |-5S_0/2\rangle_p + a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |+S_0/2\rangle_p \\ &\quad + b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |-S_0/2\rangle_p + b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |+5S_0/2\rangle_p \\ &\quad + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |+S_0/2\rangle_p + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |+7S_0/2\rangle_p) / 2. \end{aligned} \quad (17)$$

Then, we utilize the switchable cavity to achieve an appropriate Faraday rotation angle $S_0 = \pi/2$ and the photon state can also be written as

$$|\varphi\rangle_p = (e^{-i\varphi} |\sigma_{(z)}^+\rangle_p + e^{i\varphi} |\sigma_{(z)}^-\rangle_p) / \sqrt{2}, \quad (18)$$

Where

$$|\sigma_{(z)}^+\rangle_p = (|\leftrightarrow\rangle_p + i|\mathbf{b}\rangle_p) / \sqrt{2}, \quad |\sigma_{(z)}^-\rangle_p = (|\leftrightarrow\rangle_p - i|\mathbf{b}\rangle_p) / \sqrt{2}, \quad (19)$$

so the corresponding expression of the evolution state (17) could be written as

$$\begin{aligned} |\psi_{pACB}(t_{B_2})\rangle &= (a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |\square\rangle_p + a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |\square\rangle_p \\ &\quad - a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |\square\rangle_p + a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |\square\rangle_p \\ &\quad + b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |\square\rangle_p - b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |\square\rangle_p \\ &\quad + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B |\square\rangle_p + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B |\square\rangle_p) / 2 \\ &= [|\square\rangle_p (a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B + a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B \\ &\quad - b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B) \\ &\quad + |\square\rangle_p (a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B - a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B \\ &\quad + b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B)] / 2. \end{aligned} \quad (20)$$

$$\begin{aligned} |\psi_{ACB}(t_{B_2})\rangle &= (a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B + a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B \\ &\quad - b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B) / 2, \end{aligned} \quad (21)$$

$$\begin{aligned}
 |\psi'_{ACB}(t_{B_2})\rangle = & (a|\uparrow\uparrow\uparrow\rangle_A |\uparrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B - a|\uparrow\uparrow\uparrow\rangle_A |\downarrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B \\
 & + b|\downarrow\downarrow\downarrow\rangle_A |\uparrow\rangle_C |\uparrow\uparrow\uparrow\rangle_B + b|\downarrow\downarrow\downarrow\rangle_A |\downarrow\rangle_C |\downarrow\downarrow\downarrow\rangle_B) / \sqrt{2},
 \end{aligned}
 \tag{22}$$

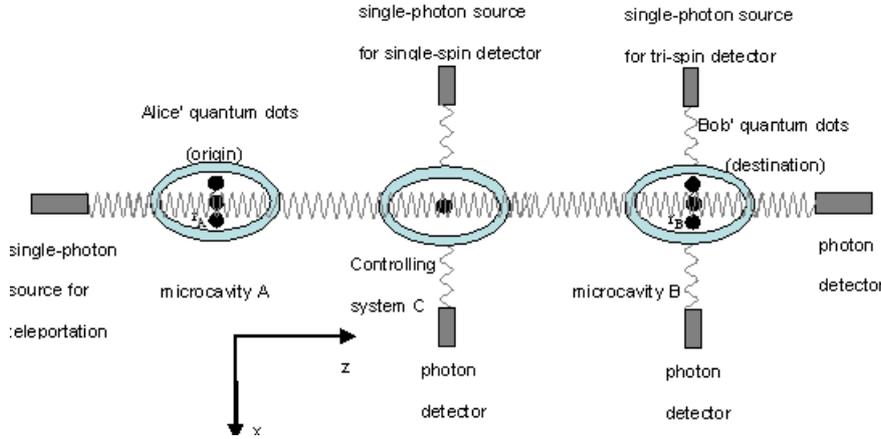


Figure 3. Controlled teleportation of a tri-qubit GHZ-Like state teleports from microcavity A to microcavity B with the help of single quantum dot at controlling system C, the photon that travels in the z direction mediates this process. Tri-spin detection and single-spin detection are performed by photon propagating along x after they entangled in microcavities.

If the linear polarization of the photon is measured first, we will obtain the following two results with equal probability.

Now, if we perform a measurement on the state of four QDs coming from microcavity A and the controlling system C, we cannot realize the CQT of tripartite GHZ-like state encoded in the electron spins of three QDs system in microcavity A to the three QDs in microcavity B through the controlling system C. Here, we implement a representation mapping for three-QD system A to complete the teleportation. In the s_x represen-

tation, there are following mappings and the total state corresponding to (20) via the above mappings becomes

$$\begin{aligned}
 |\uparrow\uparrow\uparrow\rangle = & (|\leftarrow\rangle|\leftarrow\rangle|\leftarrow\rangle + |\leftarrow\rangle|\rightarrow\rangle|\leftarrow\rangle + |\rightarrow\rangle|\leftarrow\rangle|\leftarrow\rangle \\
 & + |\rightarrow\rangle|\rightarrow\rangle|\leftarrow\rangle + |\leftarrow\rangle|\leftarrow\rangle|\rightarrow\rangle + |\leftarrow\rangle|\rightarrow\rangle|\rightarrow\rangle) / 2\sqrt{2},
 \end{aligned}
 \tag{23}$$

$$\begin{aligned}
 |\downarrow\downarrow\downarrow\rangle = & (|\leftarrow\rangle|\leftarrow\rangle|\leftarrow\rangle - |\leftarrow\rangle|\rightarrow\rangle|\leftarrow\rangle - |\rightarrow\rangle|\leftarrow\rangle|\leftarrow\rangle \\
 & + |\rightarrow\rangle|\rightarrow\rangle|\leftarrow\rangle - |\leftarrow\rangle|\leftarrow\rangle|\rightarrow\rangle + |\leftarrow\rangle|\rightarrow\rangle|\rightarrow\rangle \\
 & + |\rightarrow\rangle|\leftarrow\rangle|\rightarrow\rangle - |\rightarrow\rangle|\rightarrow\rangle|\rightarrow\rangle) / 2\sqrt{2}.
 \end{aligned}
 \tag{24}$$

$$\begin{aligned}
 |\psi_{ACB}(t_{B_2})\rangle = & \frac{1}{\sqrt{2}} \left\{ \frac{1}{\sqrt{2}} [|\leftarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) \right. \\
 & + |\rightarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) \\
 & + |\leftarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) \\
 & \left. + |\rightarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) \right] \\
 & + \frac{1}{\sqrt{2}} [|\leftarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) \\
 & + |\rightarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) \\
 & + |\leftarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) \\
 & + |\rightarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) \Big] / \sqrt{4} \\
 & + \frac{1}{\sqrt{2}} \left\{ \frac{1}{\sqrt{2}} [|\leftarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) \right. \\
 & + |\rightarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) \\
 & + |\leftarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) \\
 & \left. + |\rightarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B + \delta|\uparrow\uparrow\uparrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (a|\downarrow\downarrow\downarrow\rangle_B - \delta|\uparrow\uparrow\uparrow\rangle_B) \right] \\
 & + \frac{1}{\sqrt{2}} [|\leftarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) \\
 & + |\rightarrow\rangle_x |\leftarrow\rangle_x |\leftarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\leftarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) \\
 & + |\leftarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) + |\leftarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) \\
 & \left. + |\rightarrow\rangle_x |\leftarrow\rangle_x |\rightarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B + \delta|\downarrow\downarrow\downarrow\rangle_B) + |\rightarrow\rangle_x |\rightarrow\rangle_x |\rightarrow\rangle_x (-a|\uparrow\uparrow\uparrow\rangle_B - \delta|\downarrow\downarrow\downarrow\rangle_B) \Big] / \sqrt{4} \right\}.
 \end{aligned}
 \tag{25}$$

According to the above result, we can find if the linear polarization of the photon is measured first, then the final measurement result will depend on spin orientation of four QDs in microcavity A and the controlling system C, after implementing spin measurement on three-QD system A and single-QD system C, the whole state will collapse to one of thirty-two states with equal probability. After communicating classically the outcome of the measurement, the destination receiver operates the corresponding quantum logic gate operation, the CQT can be realized successfully and the success probability of the scheme can reach 1 if the cavities are switchable to choice the appropriate Faraday rotation angle. For example, if the linear polarization state of the photon is $|\square\rangle_r$, then the sender uses the spin detector to measure both the spin state in microcavity A and in the controlling system C. If the spin state in three-QD system A is $|\leftarrow\rangle_{A_1}|\rightarrow\rangle_{A_2}|\rightarrow\rangle_{A_3}$ and the spin state in the controlling system C is $|\uparrow\rangle_c$, the whole state will collapse to

$a|\uparrow\uparrow\uparrow\rangle_B - b|\downarrow\downarrow\downarrow\rangle_B$ at once. Subsequently, the sender informs the classical information to the receiver, the receiver performs the local gate operation $X \otimes Y \otimes X$ and he will get the state that the sender wants to transport. Here the fundamental gate operations in the microcavity are $I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$, $X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$, $Y = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$, $Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$.

Obviously, compared against more usual bipartite scheme, the tripartite CQT scheme can transport more information and expand quantum network. Similarly, the scheme in this article can be easily generalized to the CQT of multipartite GHZ-like state.

4. Conclusion

We have proposed the scheme for CQT of tripartite GHZ-like state encoded in the electron spins in QDs. A controlling system is introduced in the scheme, which participates the process of quantum teleportation as a supervisor. Without the cooperation of the controlling system, the teleportation can't be realized. With the entanglement generating through the interaction between the quantum dots in microcavities and a single photon, the CQT scheme can be realized by virtue of Faraday rotation, single photon detection and electron spin orientation measurement. The success probability of the teleportation scheme can reach 1 if the cavities are switchable to choice the appropriate Faraday rotation angle. Furthermore, we can realize the CQT of multipartite GHZ-like state with similar scheme in principle. The scheme is valuable not only for quantum security communication, large capacity quantum communication and large capacity information storage, but also for quantum distributed computation and quantum network. So far we have great challenges on quantum manipulation of photons and atoms, quantum decohe-

rence and quantum measurement, so these subjects are worth studying further in future.

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